

# Introduction to Elementary Particle Physics

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Elementary Particle Physics

Lecture 8: Esfand 13, 1397

1397-98-II

## **Bosons and Fermions in Relativistic Quantum Mechanics**

## Lecture 8

### Review of lecture 6 and some remarks:

- ▶ Schrödinger equation for a free particle in non-relativistic QM:

$$i\partial_0\psi = -\frac{\nabla^2}{2m}\psi$$

- ▶ Klein-Gordon equation for a free (massive) relativistic boson (spin 0 and electrically neutral):

$$(\square + m^2)\varphi = 0, \quad \text{with} \quad \square \equiv \partial_0^2 - \nabla^2$$

- ▶ Dirac equation for a free (massive) relativistic fermion (spin 1/2 and electrically charged):

$$[i(\gamma^0\partial_0 + \gamma^i\partial_i) - m]\psi = 0$$

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with  $\gamma$ 's satisfying the **Clifford-algebra** ( $\mu, \nu = 0, 1, 2, 3$ )

$$\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}, \quad \text{with} \quad g^{\mu\nu} = \text{diag}(1, -1, -1, -1)$$

and  $\{\gamma^5, \gamma^\mu\} = 0$

### Remark 1:

- ▶ Dirac-representation (from lecture 6):

$$\gamma^0 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}, \quad \gamma^5 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

- ▶ Weyl-representation

$$\gamma^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}, \quad \gamma^5 = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

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In general:  $i(\gamma^0 \partial_0 + \gamma^i \partial_i) \psi(\mathbf{x}, t) = 0$

- ▶ Dirac  $\gamma$  matrices:

$$\gamma^0 = \begin{pmatrix} +1 & 0 \\ 0 & -1 \end{pmatrix}_{4 \times 4}, \quad \gamma^i = \begin{pmatrix} 0 & +\sigma^i \\ -\sigma^i & 0 \end{pmatrix}_{4 \times 4}$$

In addition, we define  $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$

$$\gamma^5 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}_{4 \times 4}$$

- ▶  $\gamma^5$  matrix satisfies the following eigenvalue equation:

$$\begin{aligned} \gamma^5 \psi_R &= +\psi_R, & \text{Right-handed particles} \\ \gamma^5 \psi_L &= -\psi_L, & \text{Left-handed particles} \end{aligned}$$

- ▶  $\gamma^5$  is the generator of **chirality** (space-time transformation) !!!



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### Helicity:

- ▶ We start with Dirac equation for **massless** fermions

$$\gamma^0 \partial_0 \psi(\mathbf{x}, t) = -\gamma^i \partial_i \psi(\mathbf{x}, t) = -\boldsymbol{\gamma} \cdot \nabla \psi(\mathbf{x}, t)$$

$\times \gamma^5 \gamma^0$  and use

$$(\gamma^0)^2 = 1, \quad \text{and} \quad \gamma^5 \gamma^0 \gamma^i \equiv \Sigma^i = \begin{pmatrix} \sigma^i & 0 \\ 0 & \sigma^i \end{pmatrix}_{4 \times 4}$$

to arrive at

$$\gamma^5 \partial_0 \psi(\mathbf{x}, t) = \boldsymbol{\sigma} \cdot \nabla \psi(\mathbf{x}, t), \quad \text{or} \quad \boxed{\gamma^5 p_0 \psi = \boldsymbol{\sigma} \cdot \mathbf{p} \psi}$$

- ▶ For **massless** particles  $p_0 = E = \pm |\mathbf{p}|$
- ▶ Using  $\frac{\mathbf{p}}{p_0} = \pm \frac{\mathbf{p}}{|\mathbf{p}|} \equiv \pm \hat{\mathbf{p}}$ , we obtain  $\curvearrowright$

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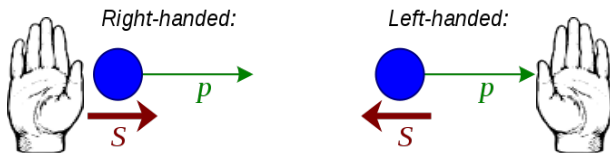
$$\mathcal{H}\psi = \pm\gamma^5\psi, \quad \text{with} \quad \mathcal{H} \equiv \boldsymbol{\sigma} \cdot \hat{\mathbf{p}}$$

**Definition:**  $\mathcal{H} \equiv \boldsymbol{\sigma} \cdot \hat{\mathbf{p}}$  is the **helicity operator**

The above relation means

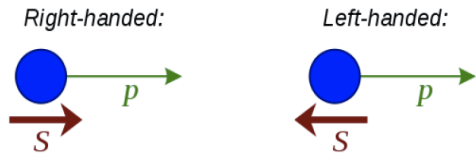
$$\text{Helicity} = \text{sgn}(E) \times \text{Chirality} \rightarrow \begin{cases} \text{Particles} & \text{sgn}(E) > 0 \\ \text{Anti-Particles} & \text{sgn}(E) < 0 \end{cases}$$

**For Particles:** Helicity=Chirality



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- Chirality/Helicity:



$$H = \frac{\vec{\sigma} \cdot \vec{p}}{|\vec{p}|}$$

Helicity = Chirality  $\times$  Sgn(E)

Chirality		Sgn(E)		Helicity
RH	+	+	Particle	+
LH	-	+		-
RH	+	-	Antiparticle	-
LH	-	-		+



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### Remark 2:

- ▶ Using Weyl-representation for  $\gamma$ -matrices, the Dirac-equation for massless fermions  $\psi = \begin{pmatrix} \chi \\ \varphi \end{pmatrix}$  reduces to

$$\begin{pmatrix} 0 & E - \boldsymbol{\sigma} \cdot \mathbf{p} \\ E + \boldsymbol{\sigma} \cdot \mathbf{p} & 0 \end{pmatrix} \begin{pmatrix} \chi \\ \varphi \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

or

$$\begin{cases} \boldsymbol{\sigma} \cdot \mathbf{p} \varphi = +E\varphi & \xrightarrow{E=\pm|\mathbf{p}|} \boldsymbol{\sigma} \cdot \hat{\mathbf{p}} \varphi = \pm\varphi, \\ \boldsymbol{\sigma} \cdot \mathbf{p} \chi = -E\chi & \xrightarrow{E=\pm|\mathbf{p}|} \boldsymbol{\sigma} \cdot \hat{\mathbf{p}} \chi = \mp\chi \end{cases}$$

- ▶ On the other hand

$$\gamma^5 \psi = \begin{pmatrix} -1 & 0 \\ 0 & +1 \end{pmatrix} \begin{pmatrix} \chi \\ \varphi \end{pmatrix} = \begin{pmatrix} -\chi \\ +\varphi \end{pmatrix},$$
$$\implies \begin{cases} \chi \text{ is a LH } \mathbf{particle} \text{ or } \mathbf{antiparticle} \\ \varphi \text{ is a RH } \mathbf{particle} \text{ or } \mathbf{antiparticle} \end{cases}$$

**Particles and antiparticles have opposite chirality**

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### Remark 3: New interpretation for $E < 0 \rightarrow$ Antiparticle

- ▶ Energy dispersion relation of a relativistic free boson or fermion

$$E^2 = \mathbf{p}^2 + m^2 \implies E = \pm \sqrt{\mathbf{p}^2 + m^2}, \quad \text{Dirac sea } \curvearrowright \gamma \rightarrow e^+ + e^-$$

- ▶ Feynman-Stueckelberg interpretation:

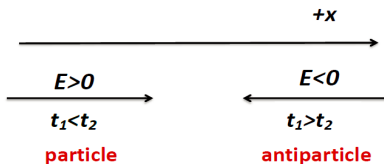
$$\psi = Ae^{iEt - i\mathbf{p}\cdot\mathbf{x}}$$

- Wave function  $\psi$  represents a **particle** with positive energy ( $+|E|$ ) and momentum  $+\mathbf{p}$  traveling in positive  $\mathbf{x}$  direction forwards in time ( $t > 0$ ).
- Same wave function  $\psi$  represents a **particle** with negative energy ( $-|E|$ ) and momentum  $-\mathbf{p}$  traveling in negative  $\mathbf{x}$  direction backwards in time ( $t < 0$ )  $\rightarrow$  An **antiparticle**

$$\psi = Ae^{i(-E)(-t) - i(-\mathbf{p})\cdot(-\mathbf{x})} = Ae^{iEt - i\mathbf{p}\cdot\mathbf{x}}$$

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### Feynman-Stueckelberg interpretation of $E < 0$



- ▶ **Pair creation** is related to conservation of electric charge (for bosons and fermions): Bosons ( $W^\pm$ ), fermions (quarks and leptons)
- ▶ The total fermion number should be conserved
- ▶ **Question:** Other conservation laws?
  - Lepton flavor number conservation
  - Quark flavor number conservation